

**Dibenzo-p-dioxin**

Twisted and puckered excited state molecular geometries

Spanget-Larsen, Jens

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## Dibenzo-*p*-dioxin. Twisted and puckered excited state molecular geometries

Jens Spanget-Larsen

*Department of Science and Environment, Roskilde University, Universitetsvej 1, DK-4000 Roskilde, Denmark.*

### *Keywords*

Dibenzo-*p*-dioxin

Excited state nuclear configurations

Density functional theory (DFT)

Time-dependent DFT

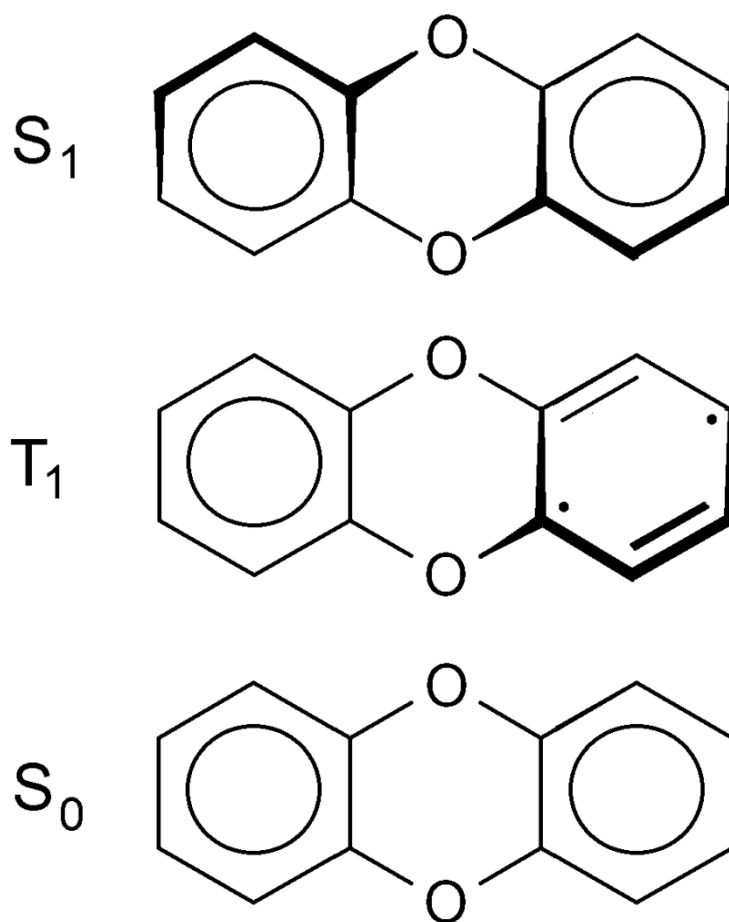
Double hybrid methods

Configuration interaction

Triplet spin populations

ABSTRACT. The title compound is generally acknowledged to have a planar  $D_{2h}$  symmetrical molecular geometry in the ground state  $S_0$ , and previous theoretical and experimental investigations seem to support the assumption of similar planar geometries in the excited singlet and triplet electronic states,  $S_1$  and  $T_1$ . But a variety of theoretical models predict non-planar equilibrium geometries for these states: In the lowest excited singlet state ( $S_1$ ) a twisted, propeller-like geometry with  $D_2$  symmetry is predicted, while a strongly puckered, biradicaloid dienediyl-like structure is predicted for the triplet state ( $T_1$ ).

*E-mail address:* Spanget@ruc.dk (J. Spanget-Larsen)



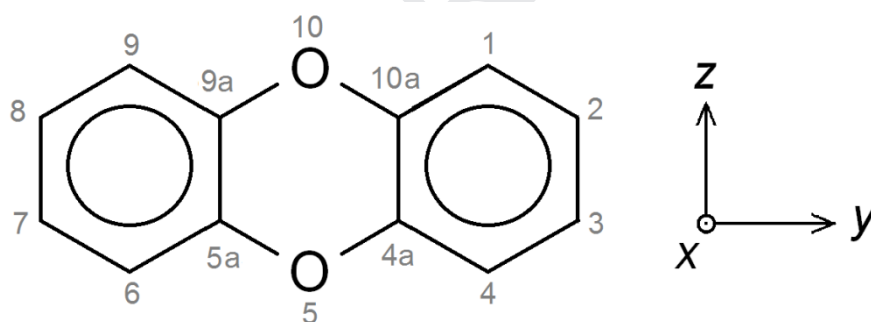
Graphical abstract

### Highlights

- Excited state molecular geometries calculated with a variety of theoretical models
- Predictions at variance with previous assumptions of planar geometries
- Twisted  $D_2$  symmetrical structure predicted for the  $S_1$  state of dibenzo-*p*-dioxin
- Puckered dienediyl-like structure predicted for the  $T_1$  state of dibenzo-*p*-dioxin

## 1. Introduction

The molecular structure of dibenzo-*p*-dioxin (DD, Scheme 1) was a matter of controversy for many years. In particular, it was discussed whether the molecule in the  $S_0$  ground state had a planar  $D_{2h}$  symmetrical structure or a non-planar “butterfly” conformation with  $C_{2v}$  symmetry (see Refs. [1-5] and literature cited therein). The question was apparently settled in 2002 when Gastilovich et al. [1] analyzed a collection of spectroscopic data and concluded that the ground state had a planar  $D_{2h}$  symmetrical nuclear configuration. Similar planar molecular geometries were deduced by Gastilovich et al. [1] for the excited electronic singlet and triplet states,  $S_1$  and  $T_1$ . These results were subsequently supported by Ljubić and Sabljčić [3] who performed *ab initio* multi-configurational CASSCF/CASPT2 calculations of several valence excited states of DD. On the basis of harmonic analyses the authors concluded that “all of the  $\pi$ - $\pi^*$  excited states exhibit planar  $D_{2h}$  minima” [3].



**Scheme 1.** Dibenzo-*p*-dioxin (DD) with atom numbering and definition of molecular coordinate system for the  $D_{2h}$  symmetrical ground state geometry.

We recently published the results of a polarization spectroscopic investigation of the electronic transitions of DD [6], leading to revision of previous assignments. It was found that time dependent density functional theory (TD-DFT) [7-9] calculations using the CAM-B3LYP [10] functional provided an adequate description of the observed vertical transitions and their polarization directions. Under the assumption of  $D_{2h}$  molecular symmetry, the lowest vertically excited singlet state  $S_1$  was predicted to be  $1^1B_{3g}$  [6], in consistency with the results of Ljubić and Sabljčić [3]. But a subsequent TD-CAM-B3LYP calculation of the

molecular equilibrium geometry for this state indicated distortion to a non-planar structure with lower symmetry. A similar result was obtained for the  $T_1$  state with CAM-B3LYP. These predictions were unexpected in view of the previously reported conclusions by Gastilovich et al. [1] and by Ljubić and Sabljčić [3]. It was thus decided to repeat the calculations with a selection of different theoretical procedures in order to check the TD–CAM-B3LYP and CAM-B3LYP predictions. The present publication reports the results of this project. Additional information is provided as Supplementary data, referred to in the ensuing text as S1-S7.

### 3. Calculations

All quantum chemical calculations were performed by using the GAUSSIAN16 software package [11]. The  $S_1$  state was computed with the CIS configuration interaction (CI) procedure [12], the TD–Hartree-Fock (TD-HF) procedure [13], and the TD-DFT [7-9] procedure with the functionals B3LYP [14], CAM-B3LYP [10], and  $\omega$ B97XD [16]. The  $T_1$  state was computed with the Hartree-Fock methods UHF and ROHF [17-19], the CI procedures CID and CISD [20-22], the DFT procedures B3LYP [14,15], CAM-B3LYP [10], M06 [23], and  $\omega$ B97XD [16], and the double hybrid procedures B2PLYP [24] and B2PLYP-D3 [25,26]. The CAM-B3LYP calculation was performed also with inclusion of a solvent contribution by using the polarizable continuum model (PCM, ethanol) [9,27]. The basis set was 6-311++G(d,p), except for CID and CISD where the basis was reduced to 6-311G(d,p) [9,11]. Geometry optimizations were performed with and without the restriction of planar  $D_{2h}$  symmetry. Attempts to optimize the  $T_1$  geometry with Møller-Plesset 2-order perturbation theory (MP2) [28,29] were unsuccessful. However, an MP2-like contribution is included in the double hybrid B2PLYP methods [24-26]. The optimized  $S_1$  and  $T_1$  geometries were characterized by harmonic analyses, except for CID, CISD, and B2PLYP-D3 where the analyses were not feasible. The computed wavefunctions were analyzed by the natural bond orbital (NBO) procedure [30], providing NBO Wiberg bond indices [31] and NBO triplet

spin populations. Selected parameters are listed in Tables 1 and 2. A number of more detailed results are given as Supplementary data, S1-S7: TD-CAM-B3LYP data for the  $S_1$  state, symmetries  $D_{2h}$  and  $D_2$  (S1-S2); CAM-B3LYP data for the  $T_1$  state, symmetries  $D_{2h}$  and  $C_1$  (S3-S4); APFD, B2PLYP, and CISD data for the  $T_1$  state, symmetry  $C_1$  (S5-S7).

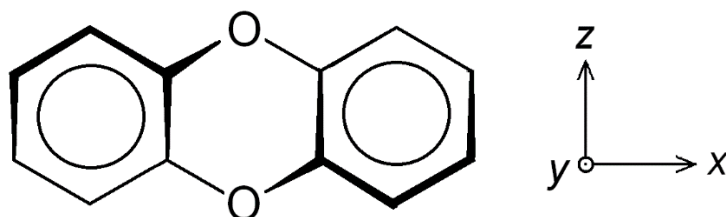
## 4. Results and Discussion

### 4.1 Singlet state $S_1$

Under the assumption of planar  $D_{2h}$  symmetrical molecular geometry, the lowest excited singlet state  $S_1$  of DD predicted by TD-CAM-B3LYP is well-described by promotion of an electron from the  $3b_{3u}(\pi)$  to the  $2a_u(\pi^*)$  molecular orbital (Fig. 1), giving rise to an  $1^1B_{3g}$  state [6] (S1). This is consistent with the result of the CASSCF/CASPT2 calculation published by Ljubić and Sabljic [3]. But a harmonic analysis of the relaxed, planar TD-CAM-B3LYP geometry reveals one imaginary frequency, corresponding to an out-of-plane twisting vibration of  $a_u$  symmetry involving the carbon nuclei of the dioxin moiety (S1). A feature at  $155\text{ cm}^{-1}$  in the phosphorescence spectrum of DD was assigned by Gastilovich et al. [1] to the corresponding  $a_u$  ground state fundamental. The prediction of an imaginary frequency for this mode indicates that the planar configuration is a transition structure. Optimization of the excited state geometry leads to a non-planar, twisted equilibrium configuration with  $D_2$  symmetry (S2). The  $1^1B_{3g}$  state is transformed to an  $1^1B_2$  state in the  $D_2$  equilibrium geometry (using the coordinate systems indicated in Schemes 1 and 2), thereby lowering the electronic energy of the system by 0.83 kcal/mol. Hence, the planar  $D_{2h}$  transition structure is interrelating two equivalent  $D_2$  symmetrical equilibrium structures with a transition energy corresponding to rapid equilibration at room temperature.

The computed molecular equilibrium geometry is shown in Fig. 2. The two carbocyclic rings are nearly planar, and bond lengths and Wiberg bond indices are insignificantly affected by the transformation from  $D_{2h}$  to  $D_2$  geometry (S1, S2). The

distortion is essentially limited to the central dioxin ring where the O-C-C-O dihedral angles are twisted by  $21^\circ$ . The twisting can be explained in part by inspection of the MO amplitudes in Fig 1:  $3b_{3u}(\pi)$  and  $2a_u(\pi^*)$  are bonding and antibonding, respectively, with respect to the C-C linkages of the dioxin moiety. Promotion of an electron from  $3b_{3u}(\pi)$  to  $2a_u(\pi^*)$  thus dilutes the  $\pi$  bonding character of these linkages, thereby reducing the tendency to planarity of the dioxin moiety in the excited state. A similar inversion distortion was predicted by Lee et al. [4] for the radical anion of DD.



**Scheme 2.** Definition of molecular coordinate system for the propeller-shaped,  $D_2$  symmetric molecular geometry predicted for the excited singlet state  $S_1$  of dibenzo-*p*-dioxin (DD).

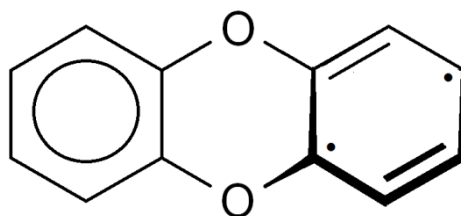
The TD-CAM-B3LYP result is at variance with previous conclusions [1-3]. The calculation was therefore repeated with a number of different theoretical procedures, see Section 3 for details. All of them predict  $D_2$  symmetrical  $S_1$  equilibrium geometries, thereby supporting the TD-CAM-B3LYP results discussed above. As listed in Table 1, the predicted O-C-C-O distortion angles range from  $14^\circ$  to  $21^\circ$ , and the energy differences  $\Delta E = E(1^1B_{3g}) - E(1^1B_2)$  between the relaxed, planar configuration ( $D_{2h}$ ) and the non-planar equilibrium geometry ( $D_2$ ) range from 0.31 to 1.00 kcal mol $^{-1}$ .

#### 4.2 Triplet state $T_1$

Under the assumption of planar  $D_{2h}$  symmetrical molecular geometry, the lowest triplet state  $T_1$  of DD predicted with CAM-B3LYP is the  $1^3B_{3g}$  state (S3), consistent with the results of Ljubić and Sabljčić [3]. But as in the case of the related singlet state  $1^1B_{3g}$ , the harmonic analysis reveals one imaginary frequency corresponding to an out-of-plane vibration of  $a_u$  symmetry (S3). Full optimization of the geometry leads to a strongly distorted,



non-planar equilibrium geometry with complete loss of symmetry, and with a biradicaloid dienediyl-like spin distribution localized in one of the carbocyclic rings, as indicated in Scheme 3 (S4). The resulting  $1^3A$  triplet state is lowered by 7.94 kcal mol<sup>-1</sup> relative to the  $1^3B_{3g}$  state of the relaxed, planar  $D_{2h}$  configuration. The redistribution of the electron density creates a dipole moment of 0.89 Debye (S4). Approximating the influence of an alcoholic solvent by means of the PCM model (solvent=ethanol) does not substantially affect the results of the calculation: The calculated molecular geometry is essentially unchanged, the energy lowering is increased from 7.94 to 8.17 kcal mol<sup>-1</sup>, and the dipole moment is increased from 0.89 to 1.21 Debye.



**Scheme 3.**

Approximate representation of the puckered, biradicaloid molecular geometry predicted for the triplet state  $T_1$  of dibenzo-*p*-dioxin (DD).

The CAM-B3LYP equilibrium geometry for the  $1^3A$  state is shown in Fig. 3. It is apparent that the two carbocyclic rings are essentially planar, while the dioxin moiety is strongly puckered. In particular, one of the O-C-C-O the dihedral angles is predicted to be as large as 42° (here taken as O5-C4a-C10a-O10). The bond lengths and Wiberg bond indices indicate that one of the benzene rings is basically intact, while the other is transformed to a dienediyl-like structure (S4). The NBO analysis reveals large spin populations on the C4a and C2 centers,  $s(C4a) = 0.74$  and  $s(C2) = 0.58$  (Fig. 4). The C4a carbon appears as a fairly well-developed, pyramidalized radical center, linked with essential single bonds to the neighbouring atoms. Atomic coordinates, NBO net charges and spin populations, dipole moment, bond lengths, and NBO Wiberg bond indices are given in S4.

The results of the CAM-B3LYP calculation are in strong disagreement with previous conclusions concerning the molecular and electronic structure of the  $T_1$  state of DD [1,3]. But

the main results of the calculation are supported by those of about a dozen different theoretical procedures, including HF, DFT, and CI methods (Section 3). All the applied procedures predict distortion to a puckered, low-symmetry structure with a dienediyl-like spin distribution, consistent with the results discussed above for CAM-B3LYP. As listed in Table 2, O5-C4a-C10a-O10 distortion angles are predicted in the range 38-51°, NBO spin populations  $s(\text{C4a})$  and  $s(\text{C2})$  in the ranges 0.66-0.88 and 0.51-0.80, respectively, and dipole moments in the range 0.36-0.89 Debye (1.21 Debye for CAM-B3LYP + PCM). However, the computed energy differences  $\Delta E = E(1^3\text{B}_{3g}) - E(1^3\text{A})$  vary widely. Detailed  $T_1$  results for CAM-B3LYP, APFD, B2LYP, and CISD are provided in S3 – S7.

## 5. Conclusions

The photophysical properties of dibenzo-*p*-dioxin (DD) and its derivatives have been investigated for many years under the explicit assumption that the  $S_0$  ground state and the excited electronic states  $S_1$  and  $T_1$  have planar  $D_{2h}$  symmetrical molecular geometries [32-38]. However, the theoretical results presented in this publication provide strong evidence for non-planar equilibrium geometries in the  $S_1$  and  $T_1$  states. Most remarkably, a strongly distorted nuclear configuration with a puckered dioxin moiety and a dienediyl-like spin distribution is predicted for the  $T_1$  state. The results are confirmed by the application of a variety of theoretical procedures, including Hartree-Fock, density functional, double hybrid, and configuration interaction procedures. The predictions are controversial, particularly for the  $T_1$  state, and should be checked by further studies. For example, the use of larger, correlation consistent basis sets should be considered, particularly in the case of the post-HF calculations. But it cannot be excluded that the interpretation of the photophysical properties of DD and its derivatives needs revision in view of the results presented in this publication.

## Competing interests

The author declares no competing interests.

## Appendix A. Supplementary data

Supplementary data to this article can be found online at . . .

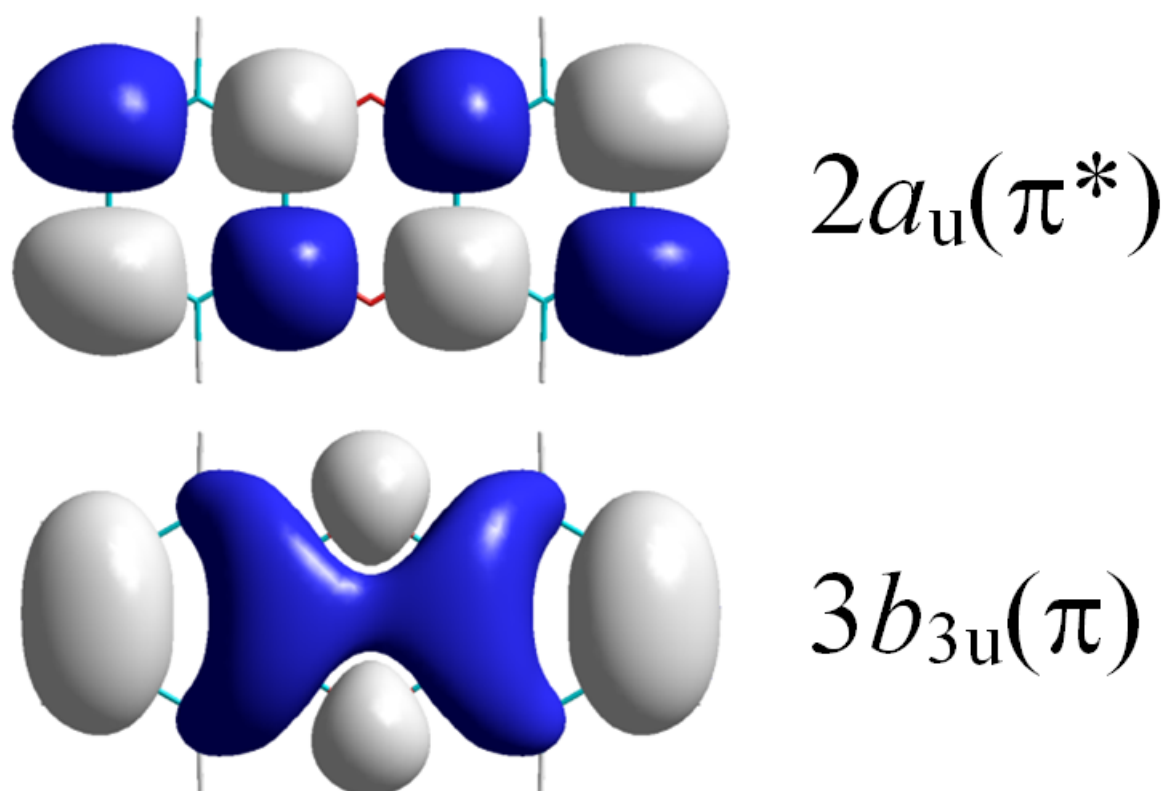
**Table 1.** Calculated singlet state  $S_1$  dihedral angles  $\Phi(\text{O-C-C-O})$  and electronic energy differences  $\Delta E = E(1^1\text{B}_{3g}) - E(1^1\text{B}_2)$  between the relaxed, planar configuration ( $D_{2h}$ ) and the non-planar equilibrium geometry ( $D_2$ ).

Method	$\Phi$	$\Delta E / \text{kcal mol}^{-1}$
TD-HF	14°	0.39
TD-B3LYP	16°	0.31
CIS	17°	1.00
TD-CAM-B3LYP	21°	0.83
TD- $\omega$ B97XD	21°	0.69

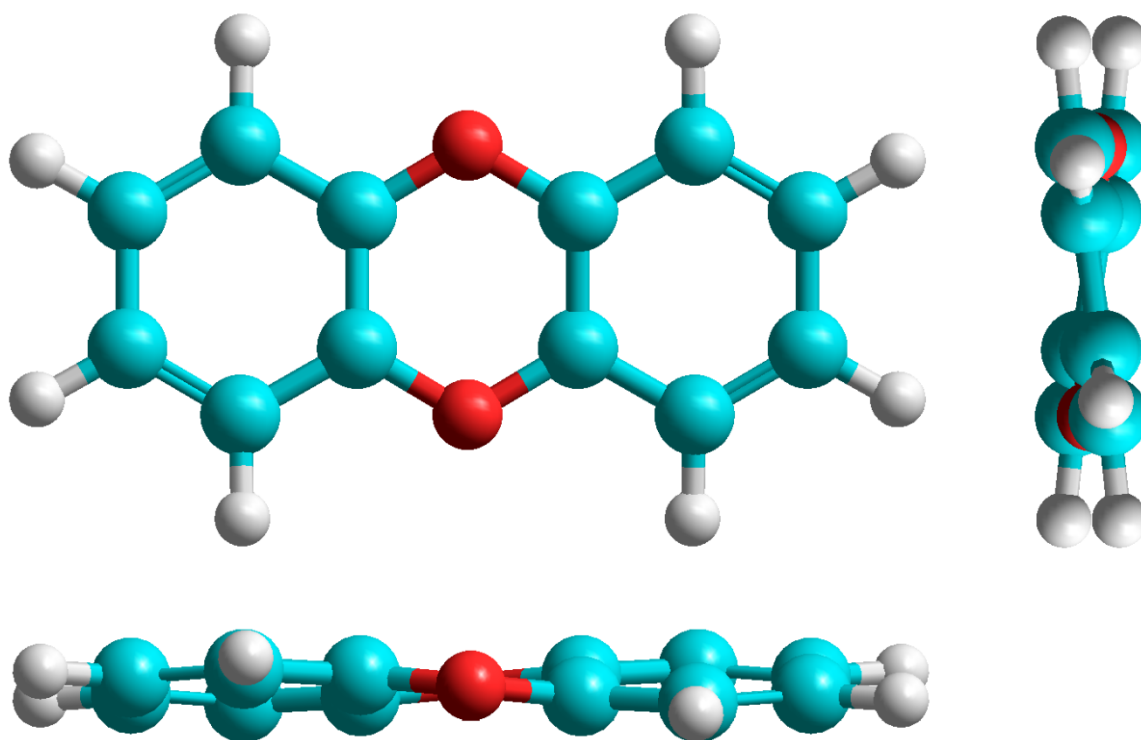
**Table 2.** Calculated triplet state  $T_1$  dihedral angles  $\Phi(\text{O5-C4a-C10a-O10})$ , NBO spin populations  $s(\text{C4a})$  and  $s(\text{C2})$ , dipole moments  $M$ , and energy differences  $\Delta E = E(1^3\text{B}_{3g}) - E(1^3\text{A})$  between the relaxed, planar configuration ( $D_{2h}$ ) and the non-planar equilibrium geometry ( $C_1$ ).

Method	$\Phi$	$s(\text{C4a})$	$s(\text{C2})$	$M / \text{Debye}$	$\Delta E / \text{kcal mol}^{-1}$
B3LYP	38°	0.66	0.51	1.15	1.60
M06	38°	0.67	0.54	0.93	1.87
APFD	38°	0.69	0.53	1.01	2.19
ROHF	39°	0.82	0.75	0.36	35.17
B2PLYP	42°	0.78	0.64	0.49	0.02
CAM-B3LYP	42°	0.74	0.58	0.89	7.94
CAM-B3LYP + PCM <sup>a</sup>	43°	0.74	0.58	1.21	8.17
B2PLYP-D3	43°	0.78	0.64	0.49	0.08
$\omega$ B97XD	43°	0.74	0.57	0.88	8.00
CISD	44°	0.88	0.80	0.39	22.03
CID	51°	0.88	0.80	0.39	21.85
UHF	51°	0.86	0.79	0.61	39.07

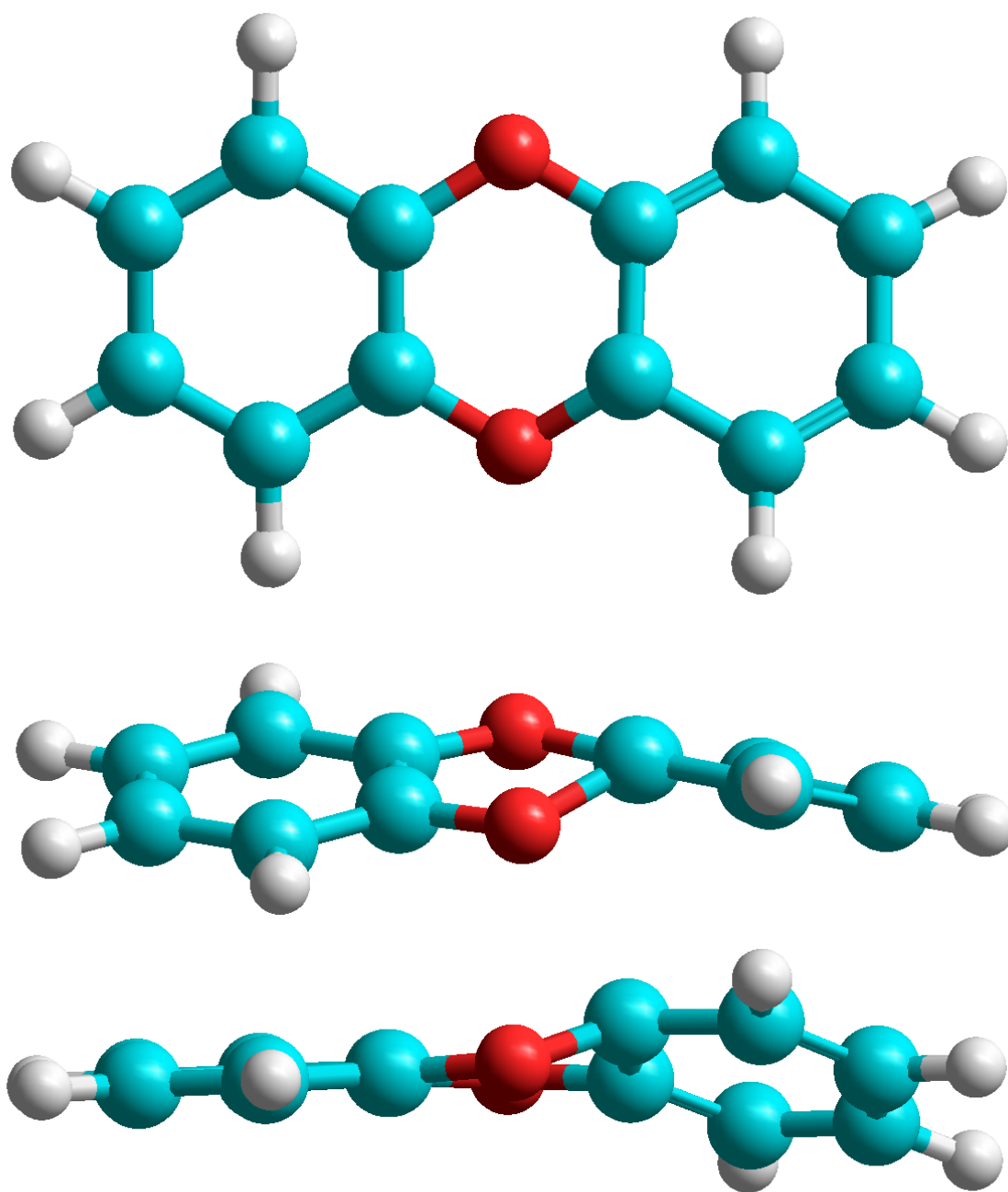
<sup>a</sup> Solvent effect included by using the PCM model, solvent=ethanol, see text.



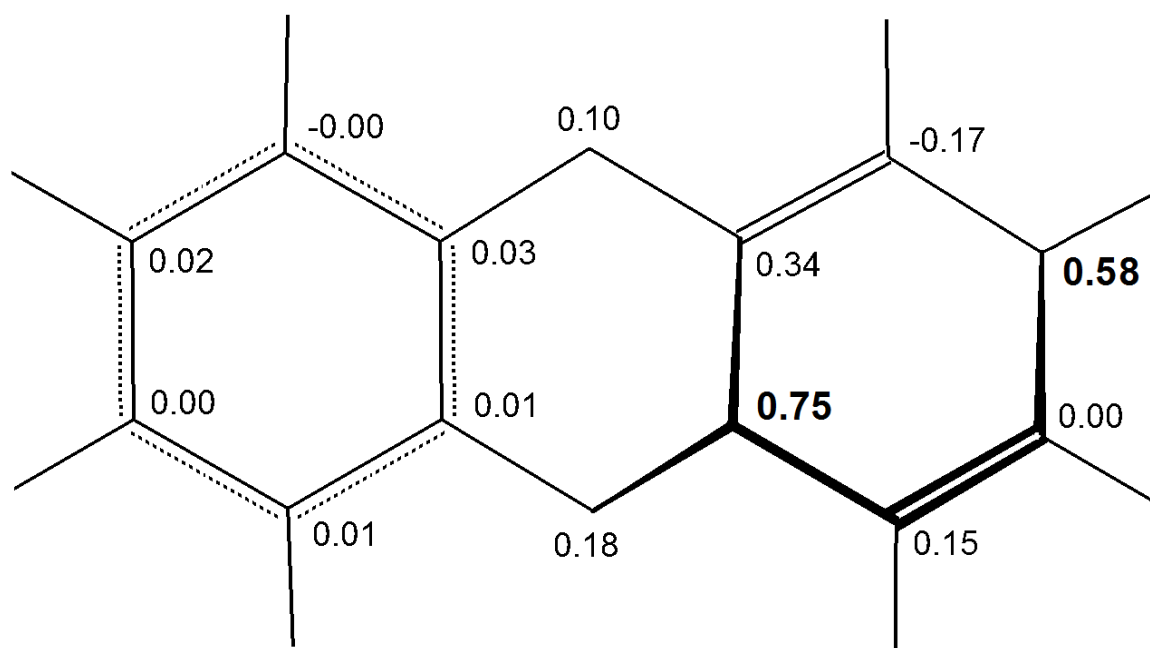
**Fig. 1.** Diagrams of the highest occupied and lowest unoccupied  $\pi$  type MOs for planar  $D_{2h}$  symmetrical dibenzo-*p*-dioxin (DD) predicted with CAM-B3LYP [6].



**Fig. 2.** Molecular equilibrium geometry of dibenzo-*p*-dioxin (DD) in the  $S_1$  singlet state ( $1^1B_2$ ) predicted with TD-CAM-B3LYP. The structure is viewed along the three different axes of symmetry in the  $D_2$  point group.



**Fig. 3.** Molecular equilibrium geometry of dibenzo-*p*-dioxin (DD) in the  $T_1$  triplet state ( $1^3A$ ) predicted with CAM-B3LYP. In the two lowest diagrams, the structure is viewed along the C4a–C10a and C5a–C9a bond axes, respectively.



**Fig 4.** NBO spin populations for the  $T_1$  triplet state ( $1^3A$ ) of dibenzo-*p*-dioxin computed with CAM-B3LYP.

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